

**Real Hilbert spaces,
 $SL(2, \mathbb{R})$
modular theory
and CFT**

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Standard real Hilbert subspaces

\mathcal{H} complex Hilbert space and $H \subset \mathcal{H}$ a real linear subspace.

Symplectic complement:

$$H' \equiv \{\xi \in \mathcal{H} : \Im(\xi, \eta) = 0 \quad \forall \eta \in H\}.$$

$H' = (iH)^\perp$ (real orthogonal complement)

$$H_1 \subset H_2 \Rightarrow H'_1 \supset H'_2 .$$

A *standard* subspace H of \mathcal{H} is a closed, real linear subspace of \mathcal{H} which is both cyclic ($\overline{H + iH} = \mathcal{H}$) and separating ($H \cap iH = \{0\}$). H is standard iff H' is standard.

H standard subspace \rightarrow anti-linear operator $S \equiv S_H : D(S) \subset \mathcal{H} \rightarrow \mathcal{H}$, where $D(S) \equiv H + iH$,

$$S : \xi + i\eta \mapsto \xi - i\eta , \quad \xi, \eta \in H .$$

$S^2 = 1 \upharpoonright_{D(S)}$. S is closed and densely defined.

Conversely, S densely defined, closed, anti-linear involution on \mathcal{H} gives

$$H = \{\xi : S\xi = \xi\} \quad \text{is a standard subspace}$$

$$H \longleftrightarrow S \quad \text{bijection}$$

Modular theory. Set

$$S_H = J_H \Delta_H^{1/2}$$

polar decomposition of $S = S_H$. Then J_H is an anti-unitary involution $\Delta \equiv S^*S > 0$

$$\Delta_H^{-it} H = H, \quad J_H H = H'$$

Borchers theorem (real subspace version)

H standard subspace, U a one-parameter group with positive generator

$$U(s)H \subset H \quad s \geq 0.$$

Then:

$$\begin{cases} \Delta_H^{it} U(s) \Delta_H^{-it} = U(e^{-2\pi t} s), \\ J_H U(s) J_H = U(-s), \end{cases} \quad t, s \in \mathbb{R}.$$

Note: Setting $K \equiv U(1)H$ we have

$$\begin{aligned}\Delta_H^{-it}K &= \Delta_H^{-it}U(1)H = U(e^{2\pi t})\Delta_H^{-it}H \\ &= U(e^{2\pi t})H \subset K, \quad t \geq 0.\end{aligned}$$

$K \subset H$ is a half-sided modular inclusion.

About the proof (adapted from Florig). With $\xi \in H, \xi' \in H'$

$$f_U(z) = (\Delta^{i\bar{z}}\xi', U(e^{2\pi z}_s)\Delta^{-iz}\xi).$$

is analytic in $\mathbb{S}_{1/2} = \{z \in \mathbb{C} : 0 < \Im z < \frac{1}{2}\}$ (the generator of $U(t)$ is positive and $\Im e^{2\pi z}_s \geq 0$ for $z \in \mathbb{S}_{1/2}$).

$V(t) = JU(-t)J$ satisfies the same assumptions then U because of $JH = H'$

$$\begin{aligned}f_U\left(t + \frac{i}{2}\right) &= (\Delta^{-1/2}\Delta^{-it}\xi', U(e^{2\pi t+i\pi}_s)\Delta^{-it}\Delta^{1/2}\xi) \\ &= (\Delta^{-1/2}\Delta^{-it}\xi', JV(e^{2\pi t}_s)\Delta^{-it}\xi) \\ &= (\Delta^{-it}\xi', (J\Delta^{1/2})V(e^{2\pi t}_s)\Delta^{-it}\xi) \\ &= (\Delta^{-it}\xi', V(e^{2\pi t}_s)\Delta^{-it}\xi) = f_V(t)\end{aligned}$$

(KMS and positivity of energy) analogously $V(t) = JU(-t)J$ satisfies the same assumptions then U because of $JH = H'$

$$f_V \left(t + \frac{i}{2} \right) = f_U(t)$$

f_U and f_V glue to an entire bounded function, thus constant.

Converse: Wiesbrock, Borchers, Araki-Zsido theorem (real subspace version)

Let H, K be standard subspaces. Assume half-sided modular inclusion:

$$\Delta_H^{-it} K \subset K, \quad t \geq 0$$

Then $\{\Delta_K^{it}, \Delta_H^{is}\}$ generates a unitary representation of the “ $ax+b$ ” group with positive energy

$$\text{dilation group} = \Delta_H^{-is/2\pi}$$

gen. of translations $P = \frac{1}{2\pi} (\log \Delta_K - \log \Delta_H)$

Conclusion:

positive energy rep. of “ $ax + b$ ” group



Borchers pair (U, H)



half-sided modular inclusion of standard subspace

therefore, if U has no non-zero fixed vector,
 (U, H) is *unique* up to multiplicity.

von Neumann algebras and real Hilbert subspaces

M von Neumann algebra on \mathcal{H} , $\Omega \in \mathcal{H}$ a cyclic separating vector,

$$H_M \equiv \overline{M_{sa}\Omega}$$

is a standard subspace of H

$$\Delta_M = \Delta_{H_M}, \quad J_M = J_{H_M}$$

In particular

$$H'_M = H_{M'}$$

Borchers theorem (original for vN algebras)

M von Neumann algebra, Ω cyclic separating vector, U a one-parameter group with positive generator with $U(s)\Omega = \Omega$ and

$$U(s)MU(-s) \subset M \quad s \geq 0.$$

Then:

$$\begin{cases} \Delta_M^{it} U(s) \Delta_M^{-it} = U(e^{-2\pi t} s), \\ J_M U(s) J_M = U(-s), \end{cases} \quad t, s \in \mathbb{R}.$$

Note: If Ω is the unique U -fixed vector then M is a type III_1 factor.

Wiesbrock, Borchers, Araki-Zsido theorem (original for vN algebras)

Let M, N be vN algebras, Ω jointly cyclic and separating vector. Assume half-sided modular inclusion:

$$\Delta_M^{-it} N \Delta_M^{it} \subset N, \quad t \geq 0.$$

Then $\{\Delta_N^{it}, \Delta_M^{is}\}$ generates a unitary representation of the “ $ax+b$ ” group with positive energy

$$\text{dilations} = \Delta_M^{-is/2\pi}$$

gen. of translations $P = \frac{1}{2\pi} (\log \Delta_N - \log \Delta_M)$

Therefore Borchers triple \Leftrightarrow Wiesbrock triple.

How many Borchers triples there are?

Is it possible that $U(s)MU(-s)' \cap M = \mathbb{C}$ for $s > 0$?

Möbius covariant nets of real Hilbert subspaces

A *local Möbius covariant net* of standard subspaces \mathcal{A} of real Hilbert subspaces on the intervals of S^1 is a map

$$I \rightarrow H(I)$$

with

1. *Isotony* : If I_1, I_2 are intervals and $I_1 \subset I_2$, then

$$H(I_1) \subset H(I_2) .$$

2. *Möbius invariance*: There is a unitary representation U of \mathbf{Mob} on \mathcal{H} such that

$$U(g)H(I) = H(gI) , \quad g \in \mathbf{Mob}, I \in \mathcal{I}.$$

Here $\mathbf{Mob} \simeq PSL(2, \mathbb{R})$ acts on S^1 as usual.

3. Positivity of the energy : $L_0 \geq 0$
4. Cyclicity : *the complex linear span of all spaces $H(I)$ is dense in \mathcal{H} .*
5. Locality : *If I_1 and I_2 are disjoint intervals then*

$$H(I_1) \subset H(I_2)'$$

First consequences

Irreducibility: $\overline{\text{real lin. span}_{I \in \mathcal{I}} \mathcal{H}(I)} = H.$

Reeh-Schlieder theorem: $H(I)$ is a standard subspace for every I .

Bisognano-Wichmann property: Tomita-Takesaki modular operator Δ_I and conjugation J_I of

$H(I)$, are

$$U(\Lambda_I(2\pi t)) = \Delta_I^{-it}, \quad t \in \mathbb{R}, \quad \text{dilations}$$

$$U(r_I) = J_I \quad \text{reflection}$$

$(\Lambda_{I_1}(t)x = e^t x, x \in \mathbb{R}, I_1 \simeq \mathbb{R}^+$ upper semi-circle)

Haag duality: $H(I)' = H(I')$ ($I' \equiv S^1 \setminus I$).

Factoriality: $H(I) \cap H(I)' = 0$

Additivity: $I \subset \cup_i I_i \implies H(I) \subset \overline{\text{real lin. span}_i H(I_i)}$.

Modular theory and representations of $SL(2, \mathbb{R})$
(Brunetti, Guido, L.)

U a unitary, positive energy representation of **Mob** on \mathcal{H} and J anti-unitary involution on \mathcal{H} s.t.

$$JU(g)J = U(rgr), \quad g \in \mathbf{Mob}$$

where $r : z \mapsto \bar{z}$ reflection on S^1 w.r.t. the upper semicircle I_1 . Then define

$$J_I \equiv U(g)JU(g)^*$$

where $g \in \mathbf{Mob}$ maps I_1 onto I .

$$\Delta_I^{it} \equiv U(\Lambda_I(-2\pi t)), \quad t \in \mathbb{R}$$

namely $-\frac{1}{2\pi} \log \Delta_I$ generator of dilations of I ,

$$S_I \equiv J_I \Delta_I^{1/2}$$

is a densely defined, antilinear, closed involution on \mathcal{H} .

$H(I)$ standard subspace associated with S_I

↓

Möbius covariant local net of real Hilbert spaces

A $\pm hsm$ factorization of real subspaces is a triple K_0, K_1, K_2 , where $\{K_i, i \in \mathbb{Z}_3\}$ is a set of

standard subspaces s.t. $K_i \subset K'_{i+1}$ is a \pm hsm inclusion.

Factorization



Local Möbius covariant net of real Hilbert spaces



Positive energy representation of $SL(2, \mathbb{R})/\{1, -1\}$

Note: Irr. positive energy rep. of $SL(2, \mathbb{R})/\{1, -1\}$ are parametrized by \mathbb{N}

Möbius covariant nets of vN algebras. A (local) *Möbius covariant net* \mathcal{A} on S^1 is a map

$$I \in \mathcal{I} \rightarrow \mathcal{A}(I) \subset B(\mathcal{H})$$

$\mathcal{I} \equiv$ family of proper intervals of S^1 , that satisfies:

A. Isotony. $I_1 \subset I_2 \implies \mathcal{A}(I_1) \subset \mathcal{A}(I_2)$

B. Locality. $I_1 \cap I_2 = \emptyset \implies [\mathcal{A}(I_1), \mathcal{A}(I_2)] = \{0\}$

C. Möbius covariance. \exists unitary rep. U of the Möbius group \mathbf{Mob} on \mathcal{H} such that

$$U(g)\mathcal{A}(I)U(g)^* = \mathcal{A}(gI), \quad g \in \mathbf{Mob}, \quad I \in \mathcal{I}.$$

D. Positivity of the energy. Generator L_0 of rotation subgroup of U (conformal Hamiltonian) is positive.

E. Existence of the vacuum. $\exists!$ U -invariant vector $\Omega \in \mathcal{H}$ (vacuum vector), and Ω is cyclic

for the von Neumann algebra $\bigvee_{I \in \mathcal{I}} \mathcal{A}(I)$ and unique U -invariant.

First consequences

Irreducibility: $\bigvee_{I \in \mathcal{I}} \mathcal{A}(I) = B(\mathcal{H})$.

Reeh-Schlieder theorem: Ω is cyclic and separating for each $\mathcal{A}(I)$.

Bisognano-Wichmann property: Tomita-Takesaki modular operator Δ_I and conjugation J_I of $(\mathcal{A}(I), \Omega)$, are

$$\begin{aligned} U(\Lambda_I(2\pi t)) &= \Delta_I^{it}, \quad t \in \mathbb{R}, && \text{dilations} \\ U(r_I) &= J_I && \text{reflection} \end{aligned}$$

(Guido-L., Frölich-Gabbiani)

Haag duality: $\mathcal{A}(I)' = \mathcal{A}(I')$

Factoriality: $\mathcal{A}(I)$ is III₁-factor (or $\mathcal{A}(I) = \mathbb{C}$).

Additivity: $I \subset \cup_i I_i \implies \mathcal{A}(I) \subset \vee_i \mathcal{A}(I_i)$ (Freedenhagen, Jorss).

- Net of factors on $\mathcal{H} \rightarrow$ Net of standard subspaces (not one-to-one) on \mathcal{H}

- Net of standard subspaces on $\mathcal{H} \rightarrow$ Net of factors on $e^{\mathcal{H}}$ (second quantization)

$$\mathcal{A}(I) \equiv \{W(h) : h \in H(I)\}''$$

Further selection properties.

• *Split property.* \mathcal{A} is *split* if the von Neumann algebra

$$\boxed{\mathcal{A}(I_1) \vee \mathcal{A}(I_2) \simeq \mathcal{A}(I_1) \otimes \mathcal{A}(I_2)}$$

(natural isomorphism) if $\bar{I}_1 \cap \bar{I}_2 = \emptyset$.

- Split is a property of the net (not of U).

- Split is crucial, e.g. for local charges, complete rationality, hyperfiniteness, classification...

• *Trace class condition.*

$$\boxed{\text{Tr}(e^{-tL_0}) < \infty, \forall t > 0}$$

- Trace class condition is standard in CFT

- Trace class condition \implies split

- Trace class condition can be refined to *log-ellipticity*

$$\log \text{Tr}(e^{-tL_0}) \sim \frac{1}{t^\alpha} (a_0 + a_1 t + \dots) \quad \text{as } t \rightarrow 0^+$$

$\alpha = 1$ (Kawahigashi, L.)

- Trace class is a property of U (not of the net).

• *Buchholz-Wichmann nuclearity:*

$$\boxed{\Phi_I^{\text{BW}}(\beta) : x \in \mathcal{A}(I) \rightarrow e^{-\beta P} x \Omega \in \mathcal{H}}$$

is nuclear, I interval of \mathbb{R} , $\beta > 0$. P translation generator (Hamiltonian).

Recall: $A : X \rightarrow Y$ is nuclear if \exists sequences $f_k \in X^*$ and $y_k \in Y$ s.t. $\sum_k \|f_k\| \|y_k\| < \infty$ and

$$Ax = \sum_k f_k(x) y_k .$$

$$\|A\|_1 \equiv \inf \sum_k \|f_k\| \|y_k\| .$$

- BW-nuclearity is a physical property (Haag-Swieca): essentially finately many localized states in a finite volume.
- BW-nuclearity is a property of the full Möbius covariant net.
- Can be refined with $\|\Phi_I^{\text{BW}}(\beta)\|_1 \leq e^{cr^m/\beta^n}$ as $\beta \rightarrow 0^+$ and \rightarrow *KMS states for translations* (Buchholz-Junglas).

Derive BW-nuclearity from the trace class condition (Buchholz, D'Antoni, L.)

- *Modular nuclearity*

M von Neumann algebra, Ω cyclic separating unit vector. Set

$$L^\infty(M) = M, \quad L^2(M) = \mathcal{H}, \quad L^1(M) = M_* .$$

Then we have the embeddings

$$\begin{array}{ccc}
 L^\infty(M) & \xrightarrow{x \rightarrow (x\Omega, J \cdot \Omega)} & L^1(M) \\
 & \searrow \Phi_{\infty,1}^M & \nearrow \\
 & \Phi_{\infty,2}^M \quad \Phi_{2,1}^M & \\
 x \rightarrow \Delta^{1/4} x\Omega & & \xi \rightarrow (\xi, J \cdot \Omega) \\
 & \searrow & \nearrow \\
 & L^2(M) &
 \end{array}$$

Now let $N \subset M$ be an inclusion of vN algebras with cyclic and separating unit vector Ω .

$L^{p,q}$ -nuclearity if $\Phi_{p,q}^M|_N$ is a nuclear operator.

$L^{\infty,2}$ -nuclearity was called *modular nuclearity*, i.e.

$$\boxed{\Phi_{\infty,2}^M|_N : x \in N \rightarrow \Delta_M^{1/4} x\Omega}$$

is nuclear.

As $\Phi_{\infty,1}^M = \Phi_{2,1}^M \Phi_{\infty,2}^M$, we have

$$\|\Phi_{\infty,1}^M|_N\|_1 \leq \|\Phi_{2,1}^M\| \cdot \|\Phi_{\infty,2}^M|_N\|_1 \leq \|\Phi_{\infty,2}^M|_N\|_1,$$

Thus

Modular nuclearity $\Rightarrow L^{\infty,1}$ – nuclearity.

indeed $\Phi_{\infty,1}^M|_N = \Phi_{2,1}^N \cdot \Phi_{\infty,2}^M|_N$ and $\|\Phi_{2,1}^N\| \leq 1$
 so $\|\Phi_{\infty,1}^M|_N\|_1 \leq \|\Phi_{\infty,2}^M|_N\|_1$. (A certain converse holds) .

- If N or M is a factor and $\Phi_{\infty,1}^M|_N$ is nuclear then $N \subset M$ is a split inclusion ($N \vee M' \simeq N \otimes M'$).

Short proof. By definition $\Phi_{\infty,1}^M|_N$ nuclear means:
 \exists sequences of elements $\varphi_k \in N^*$ and $\psi_k \in M'_*$ ($\simeq L^1(M)$) such that $\sum_k \|\varphi_k\| \|\psi_k\| < \infty$
 and

$$\omega(nm') = \sum_k \varphi_k(n)\psi_k(m') , \quad n \in N, m' \in M' .$$

where $\omega \equiv (\cdot \Omega, \Omega)$. As $\Phi_{\infty,1}^M|_N$ is normal the φ_k can be chosen normal (take the normal part). Thus the state ω on $N \odot M'$ extends to $N \otimes M'$ and this gives the split property.

Consider now the commutative diagram

$$\begin{array}{ccc}
 L^\infty(N) & \xrightarrow{\Phi_{\infty,1}^M|_N} & L^1(M) \\
 \Phi_{\infty,2}^N \downarrow & & \uparrow \Phi_{2,1}^M \\
 L^2(N) & \xrightarrow{T_{M,N} \equiv \Delta_M^{1/4} \Delta_N^{-1/4}} & L^2(M)
 \end{array}$$

$T_{M,N} \equiv \Phi_{2,2}^M|_N$. L^2 -nuclearity condition (or $L^{2,2}$ -nuclearity) means that

$$\|T_{M,N}\|_1 < \infty$$

- L^2 -nuclearity \Rightarrow modular nuclearity,

indeed $\|\Phi_{\infty,2}^M|_N\|_1 \leq \|T_{M,N}\|_1$ because $\Phi_{\infty,2}^M|_N = T_{M,N} \cdot \Phi_{\infty,2}^N$ and $\|\Phi_{\infty,2}^N\| \leq 1$.

L^2 -Nuclearity. Let $H \subset \tilde{H}$ be an inclusion of standard subspaces. Set

$$T_{\tilde{H},H} \equiv \Delta_{\tilde{H}}^{1/4} \Delta_H^{-1/4}$$

then $\|T_{\tilde{H},H}\| \leq 1$. The inclusion is *nuclear* if $T_{\tilde{H},H}$ is a nuclear (i.e. trace class) operator.

U unitary, positive energy representation of \mathbf{Mob} , $H(I)$ the associated net of standard subspaces. U satisfies L^2 *nuclearity* if $H(I) \subset H(\tilde{I})$ is nuclear if $I \subset\subset \tilde{I}$.

$SL(2, \mathbb{R})$ **identities.**

Formula 0 (Schroer-Wiesbrock)

U positive energy unitary \mathbf{Mob} rep., $\forall s \geq 0$:

$$\Delta_1^{1/4} \Delta_2^{-is} \Delta_1^{-1/4} = e^{-2\pi s L_0}$$

$\Delta_1 = \Delta_{I_1}$, $\Delta_2 = \Delta_{I_2}$, with I_1, I_2 upper and right semicircles.

About the proof. Use of double interpretation of Δ_1, Δ_2 : modular (analyticity) and $SL(2, \mathbb{R})$ (Lie algebra relations)

Formula 1 U positive energy unitary representation:

$$T_{\tilde{I},I} = e^{-sL_0} \Delta_2^{is/2\pi}$$

$s = \ell(\tilde{I}, I)$ is the inner distance (if $I = (-1, 1)$ and $\tilde{I} = (-e^s, e^s)$ on the real line, then $\ell(\tilde{I}, I) = s$) thus

$$\|T_{\tilde{I},I}\|_1 = \|e^{-sL_0}\|_1$$

About the proof.

$$\begin{aligned} e^{-2\pi s L_0} &= \Delta_1^{1/4} \Delta_2^{-is} \Delta_1^{-1/4} = \\ &\Delta_1^{1/4} \Delta_2^{-is} \left(\Delta_1^{-1/4} \Delta_2^{is} \right) \Delta_2^{-is} = T_{I_1, I_1, s} \Delta_2^{-is} \end{aligned}$$

Formula 2

$$T_{I, I_{a',a}} = e^{-a'P'_I} e^{-aP_I} e^{-iaP_I} e^{ia'P'_I} .$$

$I_{a',a} \equiv \tau'_{-a'} \tau_a I$ with $a, a' > 0$.

$$e^{-2sL_0} = e^{-\tanh(\frac{s}{2})P} e^{-\sinh(s)P'} e^{-\tanh(\frac{s}{2})P}$$

therefore

$$e^{-2sL_0} \leq e^{-2 \tanh(\frac{s}{2})P}$$

in particular $e^{-i\pi L_0} = e^{iP} e^{iP'} e^{iP}$.

About the proof. Consider $\tilde{I} = (0, \infty)$, $I = (t, \infty)$, then

$$\begin{aligned} T_{\tilde{I}, I} &= \Delta_{\tilde{I}}^{1/4} \Delta_I^{-1/4} \\ &= \left(\Delta_{\tilde{I}}^{1/4} U(t) \Delta_{\tilde{I}}^{-1/4} \right) U(-t) \\ &= e^{-tP} e^{itP} \end{aligned}$$

where we have used the Borchers commutation relation $\Delta_{\tilde{I}}^{is} e^{itP} \Delta_{\tilde{I}}^{-is} = e^{i(e^{-2\pi s})tP}$. Any $I \subset\subset \tilde{I}$ is obtain by iteration the above, get a formula and compare with formula 1.

Formula 3

$$\|e^{-\tan(2\pi\lambda)d_I P} \Delta_I^{-\lambda}\| \leq 1, \quad 0 < \lambda < 1/4.$$

with d_I the usual lenght. Thus

$$e^{-2 \tan(2\pi\lambda)d_I P} \leq \Delta_I^{2\lambda}.$$

so we have

$$e^{-2d_I P} \leq \Delta_I^{1/4} \leq e^{\frac{2}{d_I} P'}$$

Modular nuclearity and L^2 -nuclearity

L^2 -nuclearity implies modular nuclearity and $\|\Delta_{\tilde{H}}^{1/4} E_H\|_1 \leq \|T_{\tilde{H}, H}\|_1$.

Comparison of nuclearity conditions

Let H be a Möbius covariant net of real Hilbert subspaces of a Hilbert space \mathcal{H} . Consider the following nuclearity conditions for H .

Trace class condition: $\text{Tr}(e^{-sL_0}) < \infty, s > 0$;

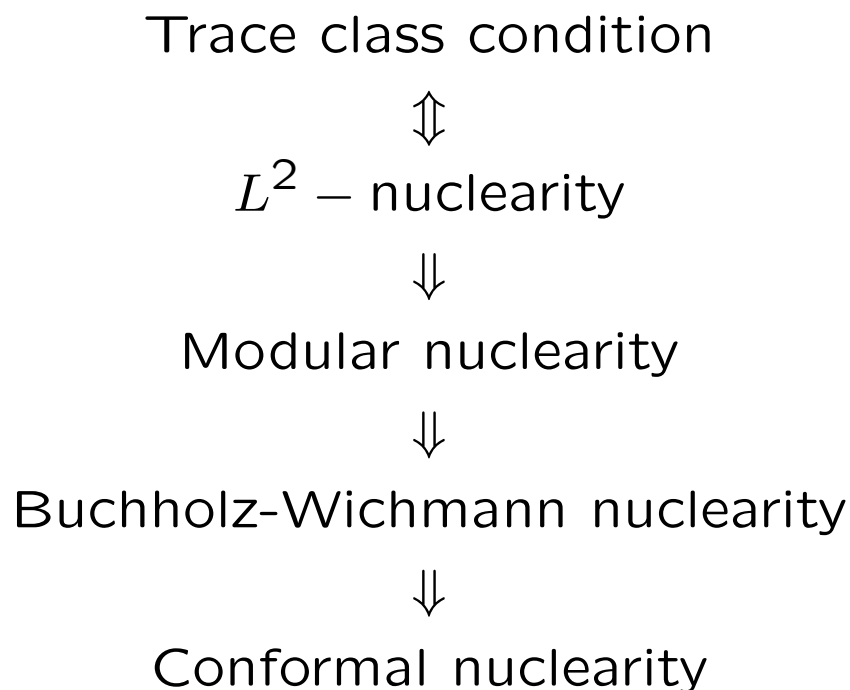
L^2 -nuclearity: $\|T_{\tilde{I}, I}\|_1 < \infty, \forall I \subset\subset \tilde{I}$;

Modular nuclearity: $\Xi_{\tilde{I}, I} : \xi \in H(I) \rightarrow \Delta_{\tilde{I}}^{1/4} \xi \in \mathcal{H}$ is nuclear $\forall I \subset\subset \tilde{I}$;

Buchholz-Wichmann nuclearity: $\Phi_I^{\text{BW}}(s) : \xi \in H(I) \rightarrow e^{-sP}\xi \in \mathcal{H}$ is nuclear, I interval of \mathbb{R} , $s > 0$ (P the generator of translations);

Conformal nuclearity: $\Psi_I(s) : \xi \in H(I) \rightarrow e^{-sL_0}\xi \in \mathcal{H}$ is nuclear, I interval of S^1 , $s > 0$.

We shall show the following chain of implications:



Where all the conditions can be understood for a specific value of the parameter, that will be

determined, or for all values in the parameter range.

We have already discussed the implications “Trace class condition $\Leftrightarrow L^2$ -nuclearity \Rightarrow Modular nuclearity”.

Modular nuclearity \Rightarrow BW-nuclearity

We have

$$\|\Phi_{I_0}^{\text{BW}}(d_I)\|_1 \leq \|\Xi_{I, I_0}\|_1$$

where d_I is the length of I on \mathbb{R} .

BW-nuclearity \Rightarrow Conformal nuclearity

By formula 2 there exists a bounded operator B with norm $\|B\| \leq 1$ such that $e^{-sL_0} = B e^{-\tanh(\frac{s}{2})H}$, therefore

$$\Psi_I(s) = B \Phi_I^{\text{BW}}(\tanh(s/2))$$

$$\|\Psi_I(s)\|_1 \leq \|\Phi_I^{\text{BW}}(\tanh(s/2))\|_1.$$

Consequences

- *Distal split property.* If $\text{Tr}(e^{-sL_0}) < \infty$ for a fixed $s > 0$, then $\mathcal{A}(I) \subset \mathcal{A}(\tilde{I})$ is split if $I \subset \tilde{I}$ and $\ell(\tilde{I}, I) > s$ e.g. free probability nets (D'Antoni, Radulescu, L.).

- *Constructing KMS states.* $\mathcal{A}|_{\mathbb{R}}$ restriction of \mathcal{A} to $\mathbb{R} \simeq S^1 \setminus \{-1\}$, \mathcal{A}_0 the quasi-local C^* -algebra. i.e. the norm closure of $\cup_I \mathcal{A}(I)$ as I varies in the bounded intervals of \mathbb{R} . Let $\mathfrak{A} \subset \mathcal{A}_0$ the C^* -algebras of elements with norm continuous orbit, namely

$$\mathfrak{A} = \{X \in \mathcal{A}_0 : \lim_{t \rightarrow 0} \|\tau_t(X) - X\| = 0\}$$

τ translation automorphism group.

Thm. If the trace class condition holds for \mathcal{A} with the asymptotic bound

$$\text{Tr}(e^{-sL_0}) \leq e^{\text{const.} \frac{1}{s^\alpha}}, \quad s \rightarrow 0^+$$

for some $\alpha > 0$, then the BW-nuclearity holds with $m = n = \alpha$.

If the trace class condition holds with log-ellipticity (above asymptotics) then for every $\beta > 0$ there exists a translation β -KMS state on \mathfrak{A} .

- *L^2 -Nuclearity and KMS states in higher dimensions.*

\mathcal{O} a double cone in the Minkowski spacetime \mathbb{R}^{d+1} , $\mathcal{A}(\mathcal{O})$ the local von Neumann algebra associated with \mathcal{O} by the $d+1$ -dimensional scalar, massless, free field.

With I an interval of the time-axis $\{x = \langle x_0, \mathbf{x} \rangle : \mathbf{x} = 0\}$ we set

$$\mathcal{A}_0(I) \equiv \mathcal{A}(\mathcal{O}_I)$$

where \mathcal{O}_I is the double cone $I'' \subset \mathbb{R}^{d+1}$, the causal envelope of I . Then \mathcal{A}_0 is a translation-dilation covariant net on \mathbb{R} . \mathcal{A}_0 is local if d is

odd and twisted local if d is even. Moreover \mathcal{A}_0 extends to a Möbius covariant net on S^1 (d odd) as one has a natural factorization.

We have:

$$\mathcal{A}_0 = \bigotimes_{k=0}^{\infty} N_d(k) \mathcal{A}^{(k)}$$

where $\mathcal{A}^{(k)}$ is the Möbius covariant net on S^1 associated with the k^{th} -derivative of the $U(1)$ -current algebra and $N_d(k)$ is a multiplicity factor (see below).

This follows because the one-particle **Mob** representation U_0 decomposes

$$U_0 = \bigoplus_{k=1}^{\infty} N_d(k) U^{(k)}$$

where $U^{(k)}$ is the positive energy irreducible representation of $PSL(2, \mathbb{R})$ with lowest weight k .

A spherical harmonics computations determines the multiplicity factor $N_d(k)$. As $k \rightarrow \infty$:

$$\begin{aligned} N_d(k+1) &= \dim(\mathcal{P}_k \ominus \mathcal{P}_{k-2}) \\ &= m_{d-1}(k-1) + m_{d-1}(k) \sim \frac{2}{(d-2)!} k^{d-2}, \end{aligned}$$

with $\mathcal{P}_k \ominus \mathcal{P}_{k-2}$ the k -spherical harmonics and $m_d(k) \sim \frac{1}{(d-1)!} k^{d-1}$. Thus

$$\log \text{Tr}(e^{-sL_0}) \sim \frac{2}{s^d} \quad s \rightarrow 0^+,$$

where L_0 is the conformal Hamiltonian of \mathcal{A}_0 .

Problems.

- $\text{Tr}(e^{-sL_0}) < \infty \Leftrightarrow$ split property?
- e^{-sL_0} compact \Leftrightarrow split property?
- $\text{Tr}(e^{-sL_0}) < \infty \Rightarrow \text{Tr}(e^{-sL_{0,\rho}}) < \infty$ in every irreducible representation ρ of \mathcal{A} ?